

# *The Future is Flavourful*

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## **Probing Axion-like Particles: Novel Detection Channels and Expanding Experimental Sensitivity for Dark Matter and Solar ALPs**

*Reference : PHYS. REV. D 108, 043029 (2023)*

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- ◆ **Interactions**
- ◆ **Dark Matter ALPs**
- ◆ **Solar ALPs**
- ◆ **Analysis and Results**
- ◆ **Summary & Future Prospects**

# Axions & ALPs : Portal to New Physics

- ✓ Some phenomena cannot be explained with current SM:
  - a. Strong CP problem
  - b. The existence of DM
- ✓ Axion - first introduced in 1970s as a solution to the strong CP problem in QCD
- ✓ Axionlike Particles (ALPs) – Variants of QCD axions – not necessarily solutions to the strong CP problem
- ✓ For QCD axions, axion-photon coupling  $\sim 1/f_a$ ,  $m_a \sim 1/f_a$ .  
For ALPs,  $m_{\text{ALP}} \sim \Lambda^2 / f_a$  is nearly arbitrary.
- ✓ Sources of ALPs:
  - (a) Dark Matter
  - (b) Sun

*Coupling of ALPs to photons ( $g_{agg}$ ) - chances to explore and understand the physics beyond SM*

# Interactions

Lagrangian:

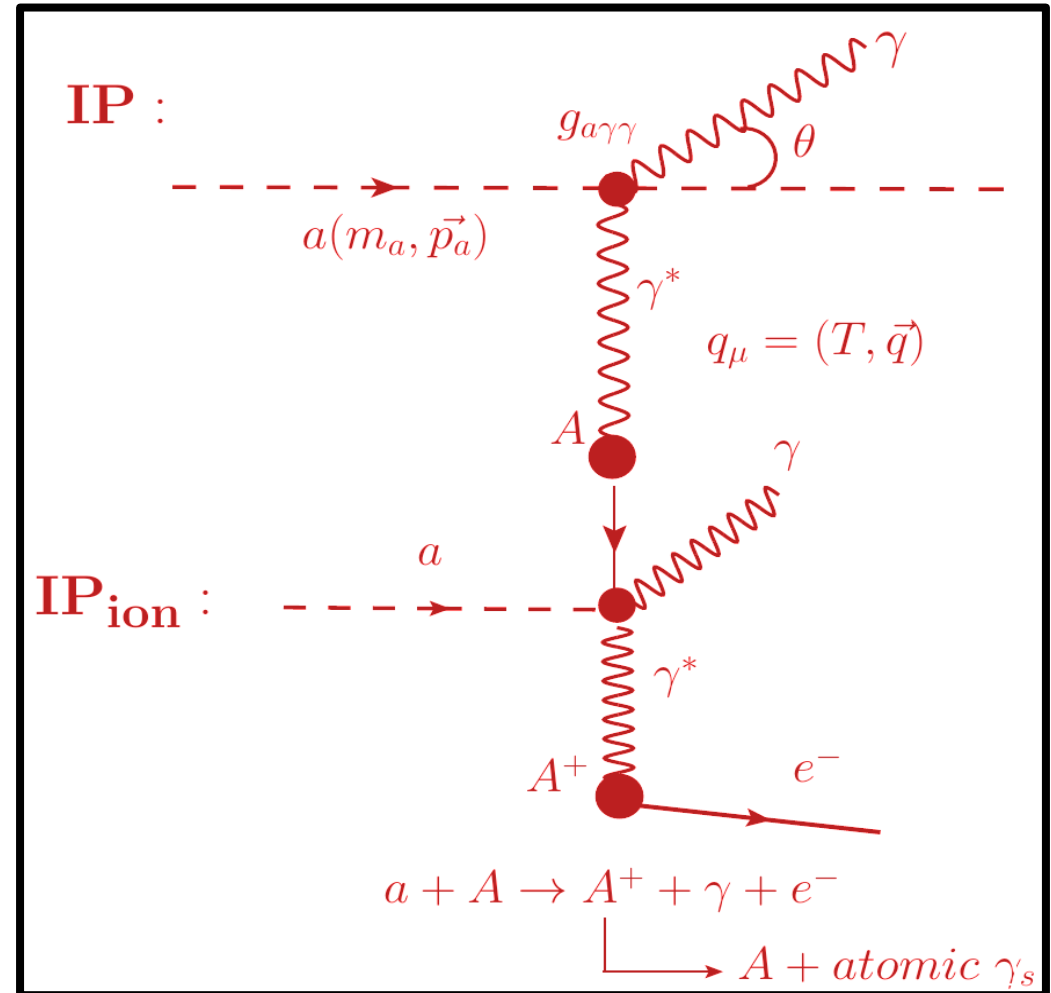
$$\mathcal{L}_I = -\frac{g_{a\gamma\gamma}}{4}\phi_a F_{\mu\nu}\tilde{F}^{\mu\nu} - \sum_f \frac{g_{aff}}{2m_f}\partial_\mu\phi_a\bar{\Psi}_f(\gamma^\mu\gamma_5)\Psi_f$$

Vacuum Decay,

$$a \rightarrow \gamma_1 + \gamma_2$$

Three Inverse Primakoff (IP) interactions are:

$$a+A \rightarrow \begin{cases} \gamma + A & \text{IP}_{el} : \text{elastic scattering} \\ \gamma + A^* & \text{IP}_{ex} : \text{atomic excitation} \\ \gamma + A^+ + e^- & \text{IP}_{ion} : \text{atomic ionization} \end{cases}$$



# Double-pole Enhancement for IP<sub>ion</sub>

Differential cross section of the ALP IP processes:

$$\frac{d^2\sigma}{dTd\Omega} = \frac{\alpha_{em}g_{a\gamma\gamma}^2}{16\pi} \frac{E_a - T}{v_a E_a} \left( \frac{V_L}{(q^2)^2} \mathcal{R}_L + \frac{V_T}{(Q^2)^2} \mathcal{R}_T \right)$$

$$\mathcal{R}_{L/T} = \sum \overline{\sum} |\langle F | \rho / \vec{j}_\perp | I \rangle|^2 \delta(E_I - E_F - T)$$

$$V_L = 2 [E_a^2 - m_a^2 + (E_a - T)^2] q^2 - (q^2)^2 - (T^2 - 2E_a T + m_a^2)^2,$$

$$V_T = m_a^4 + \frac{Q^2}{2q^2} [(m_a^4 - 4m_a^2 E_a T) + (2m_a^2 + 4E_a^2 - 4E_a T + 2T^2) Q^2 - (Q^2)^2]$$

**Exhibit a double pole structure at  $Q^2 = 0$  and for  $m_a \neq 0$**

For regulating the photon pole:

$$\frac{1}{(Q^2)^2} \rightarrow \frac{1}{(Q^2 - \Lambda_T^2/4)^2 + T^2 \Lambda_T^2} \quad \Lambda_T \equiv n_A \sigma_\gamma(T)$$

# Double-pole Enhancement for $IP_{ion}$

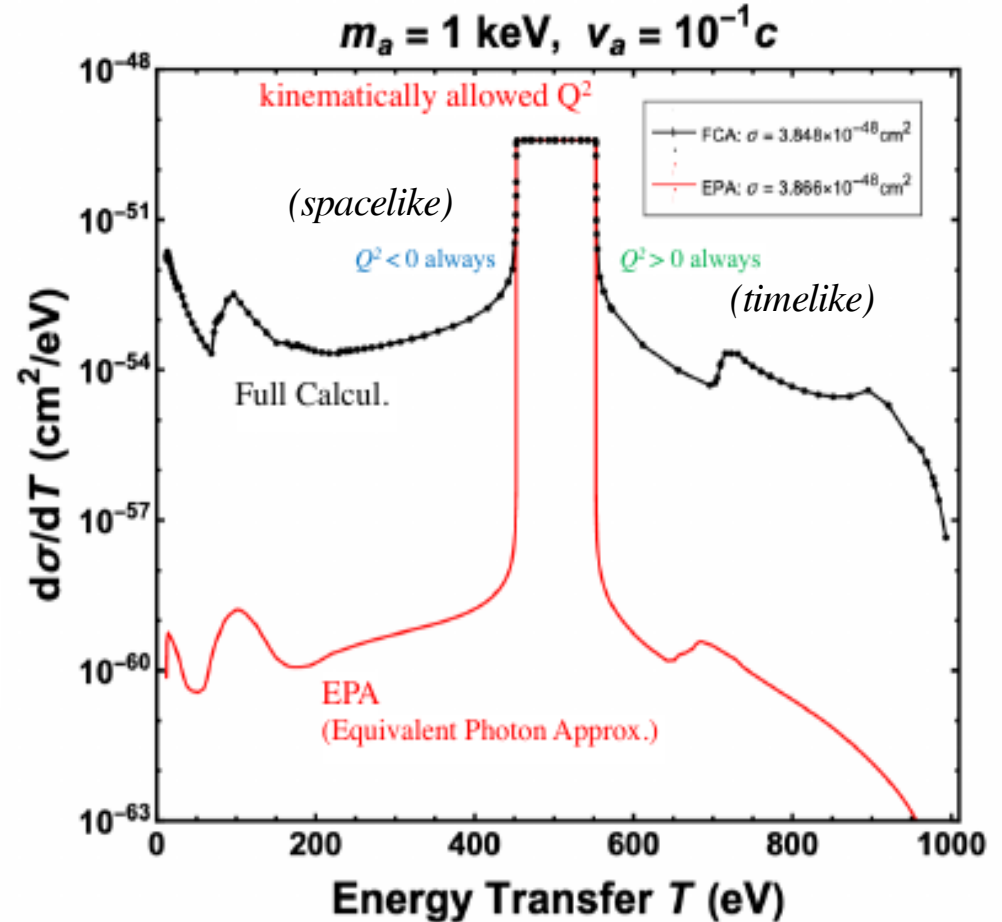
From kinematics:

$$Q^2 = m_a^2 - 2E_a(E_a - T)(1 - v_a \cos \theta)$$

◆ For Non-Relativistic ALP ( $E_a \approx m_a$ ),  $IP_{ion}$  has double pole enhancement near  $T \approx m_a/2$

➤ EPA is a good approximation in  $Q^2 = 0$  region

$$EPA : \mathcal{R}_L = 0, \mathcal{R}_T = \frac{T}{2\pi^2\alpha} \sigma_\gamma(T)$$



# Dark Matter ALPs

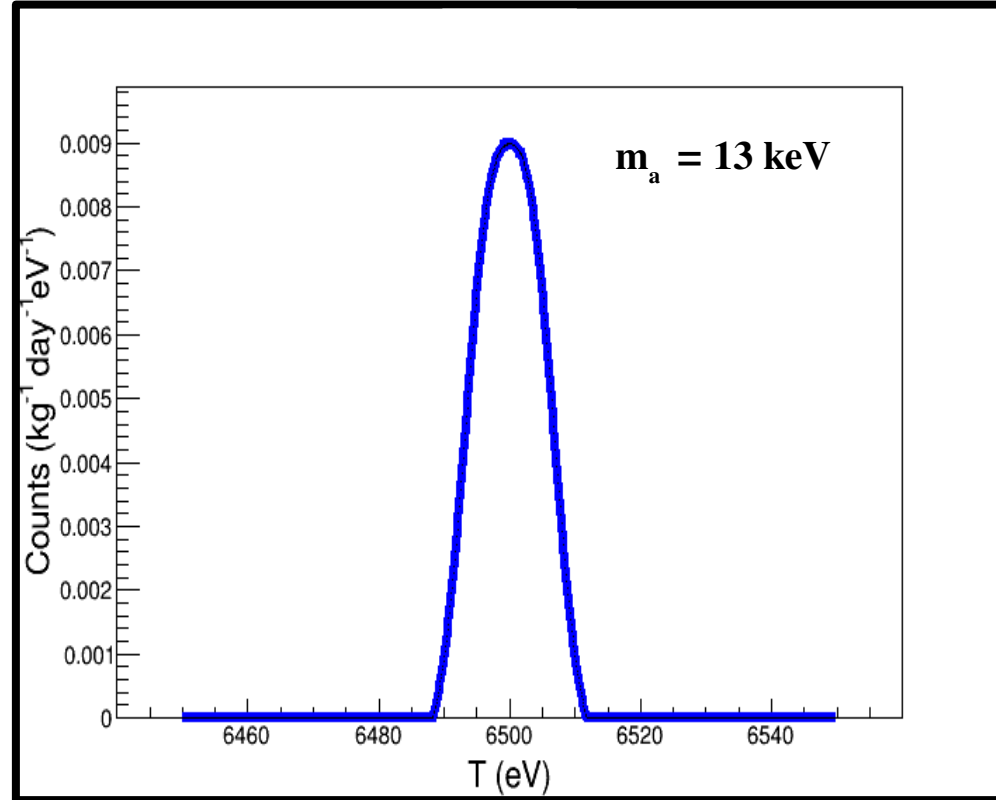
## Inverse Primakoff Ionization (IP<sub>ion</sub>)

Differential Cross-section,

$$\left. \frac{d\sigma}{dT} \right|_{EPA} = \frac{g_{a\gamma\gamma}^2}{8} \left[ \frac{1}{E_A^2 - M_A^2} \right] \frac{M_A^4}{4\pi n}$$

## Vacuum decay rate

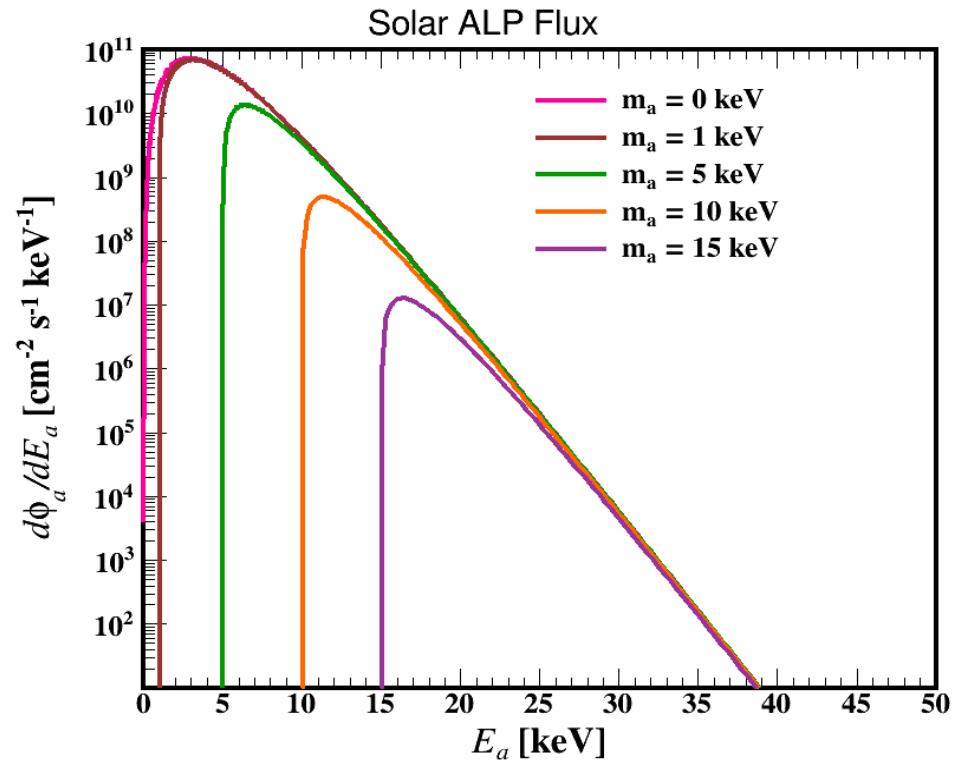
$$\Gamma_{a\gamma\gamma}^V = \frac{1}{64\pi} g_{a\gamma\gamma}^2 m_a^3$$



# Solar ALPs

- Relatively of low mass - can travel in relativistic speed
- Dominant interaction channel -  $\text{IP}_{\text{el}}$

Flux of solar axions:





# Analysis and Results

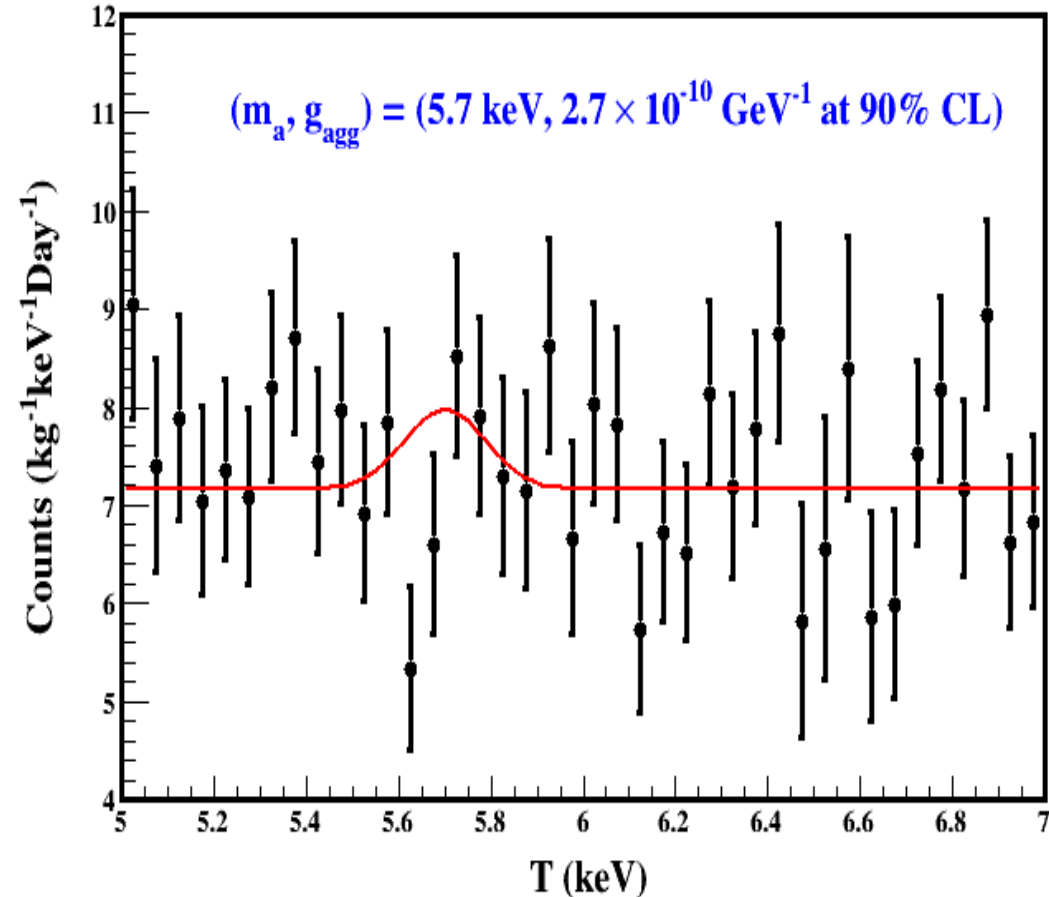
## TEXONO

(a) Low energy data with point-contact germanium detector at **300 eV** to  $12 \text{ keV}_{ee}$

(b) High-Purity Germanium detector data at  $12 \text{ keV}_{ee}$  to  $3000 \text{ keV}_{ee}$   
( $1.98 \text{ keV}_{ee}$  at  $1 \text{ MeV}_{ee}$ )

- **low threshold & excellent energy resolution**

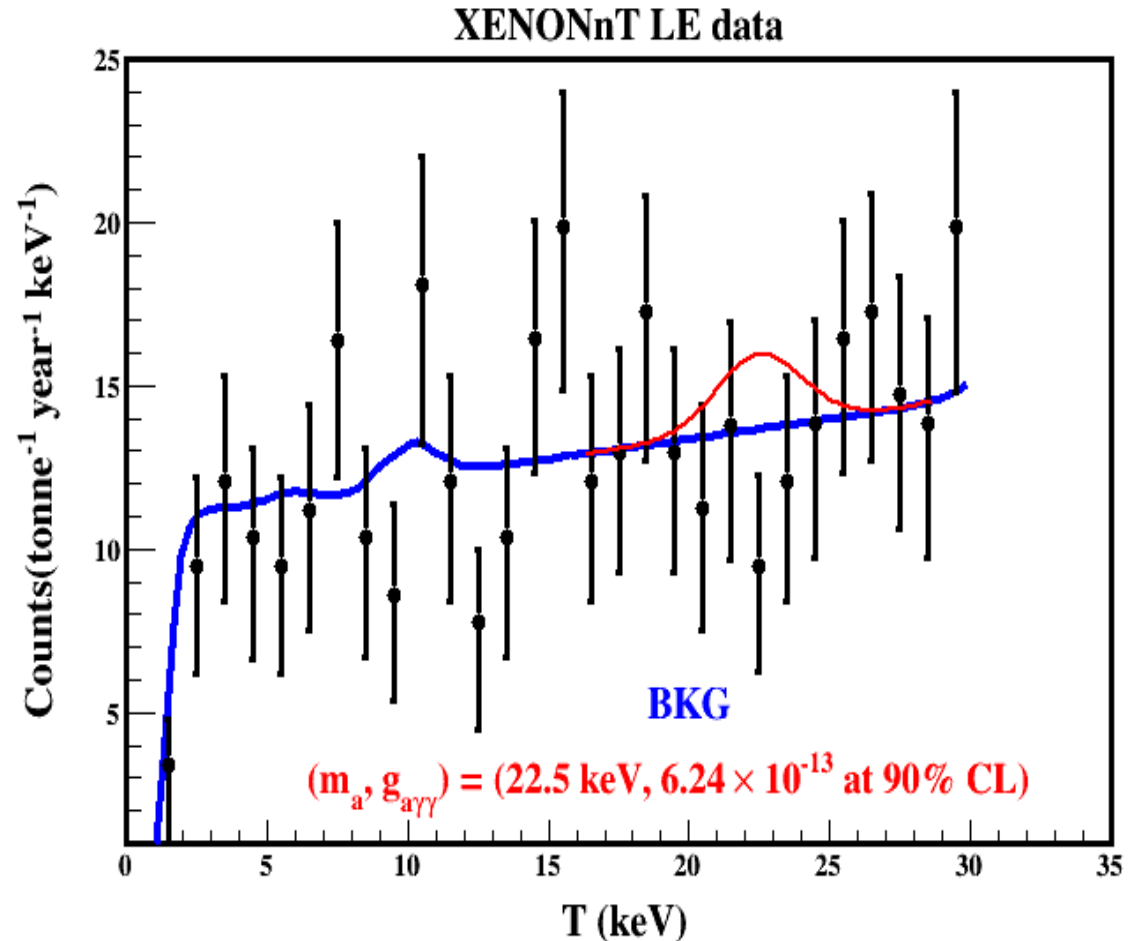
TEXONO NPC data



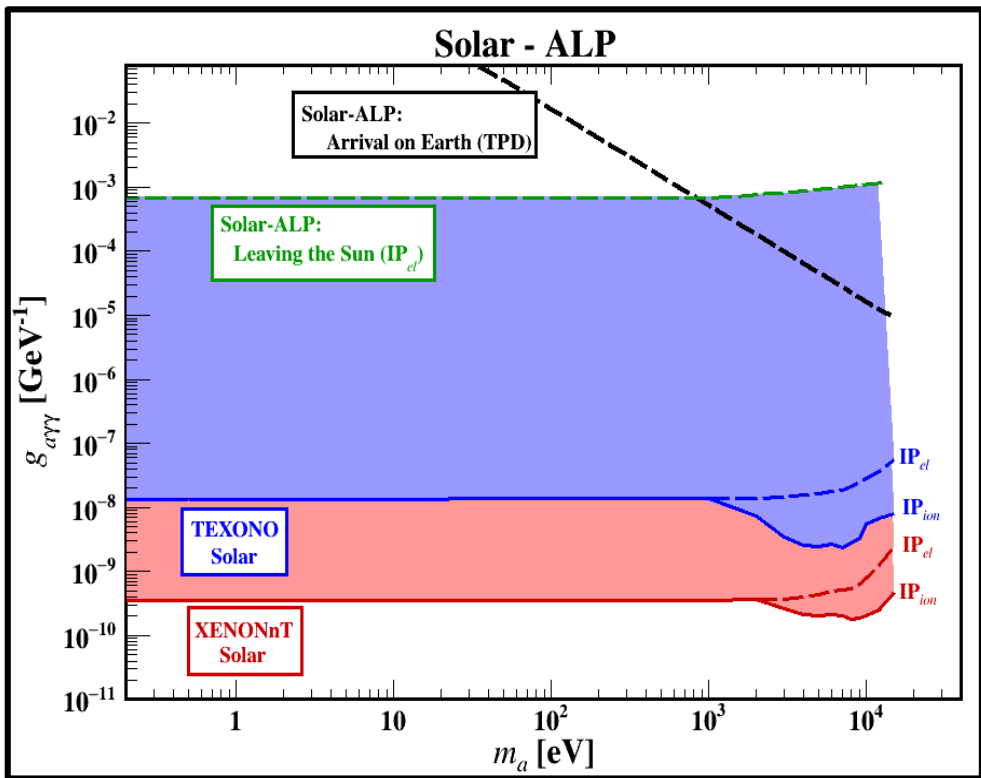
# Analysis and Results

## XENONnT

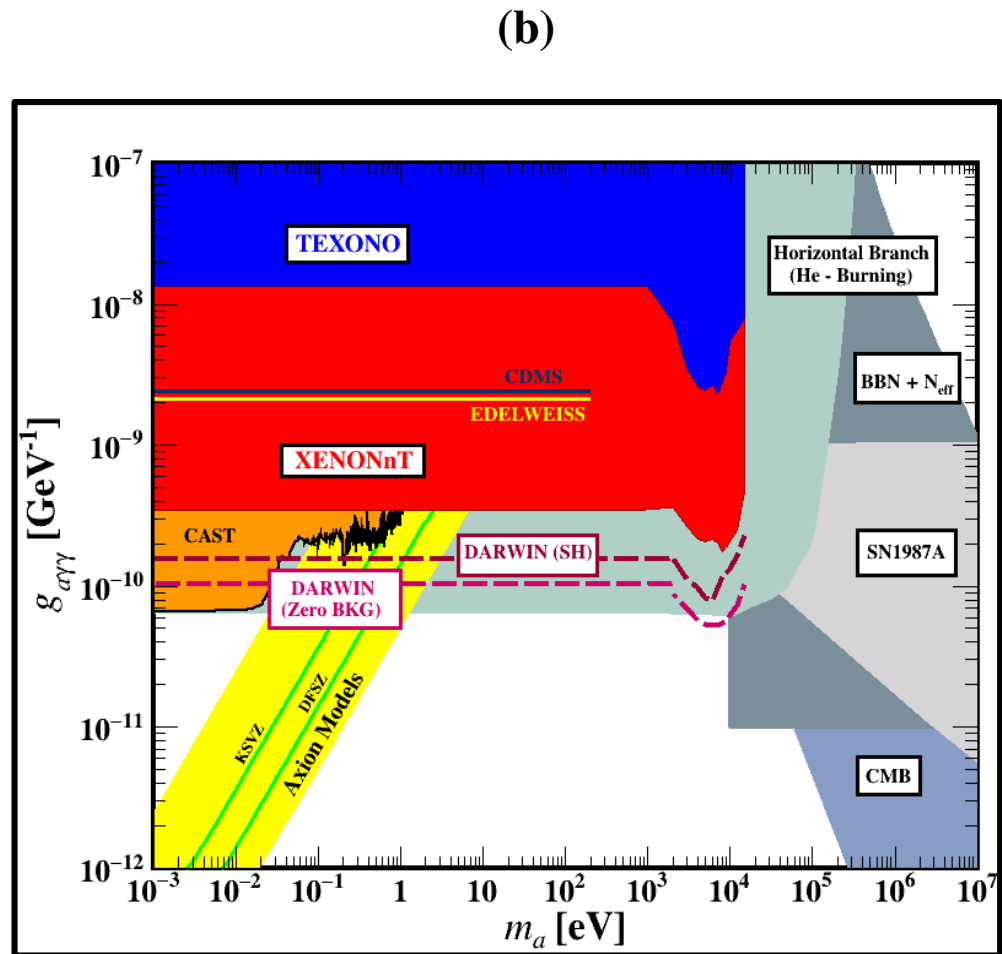
XENONnT data with liquid xenon at 1 keV<sub>ee</sub> to 140 keV<sub>ee</sub> - **large exposure** while having low threshold and background.



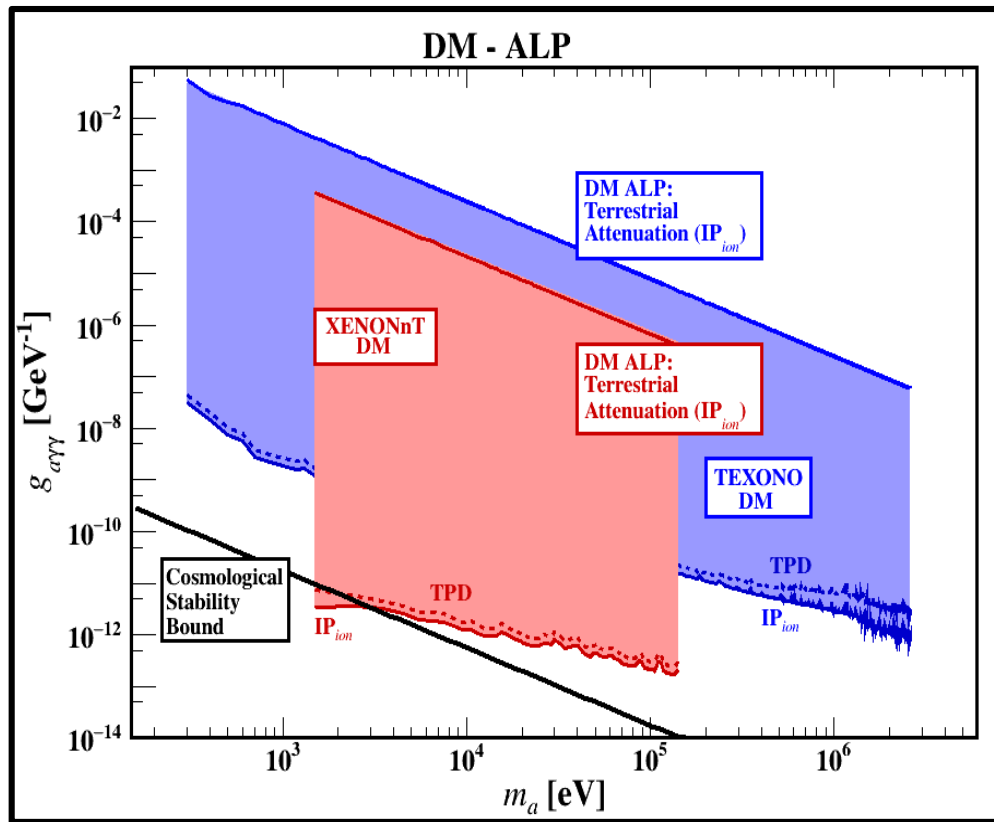
# Exclusion Plots – Solar ALP



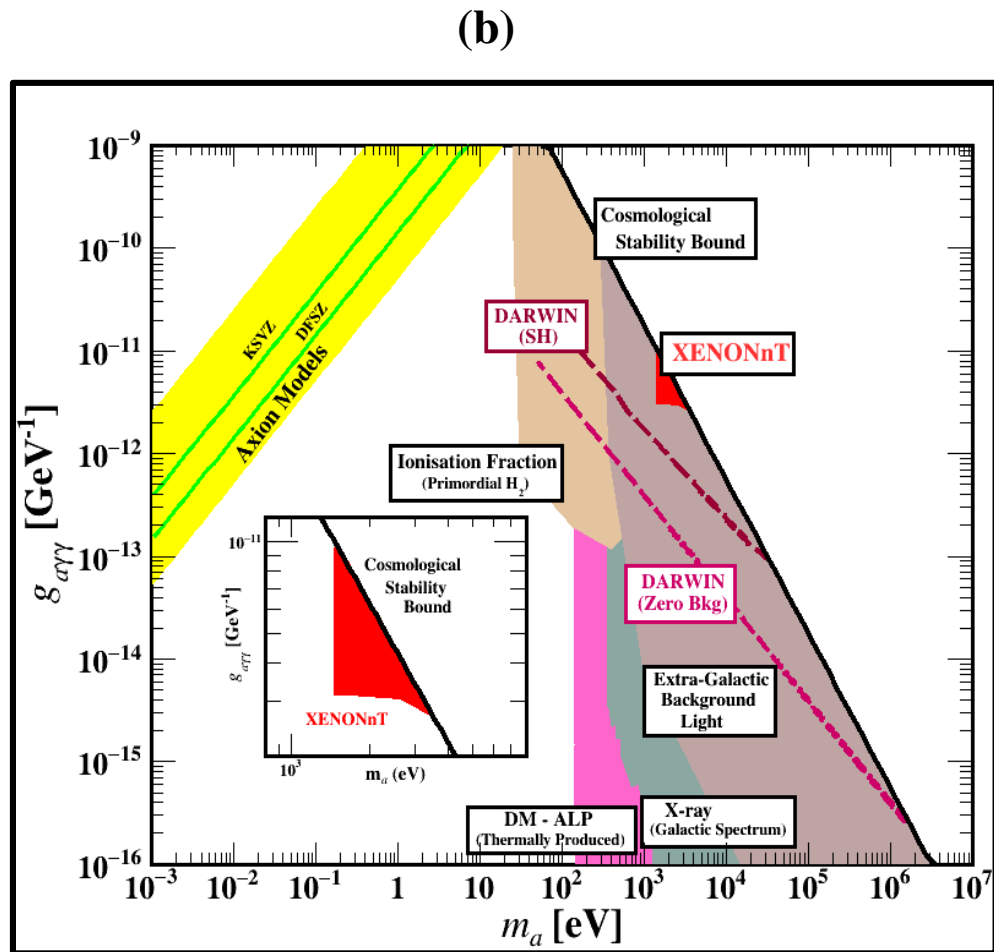
(a)



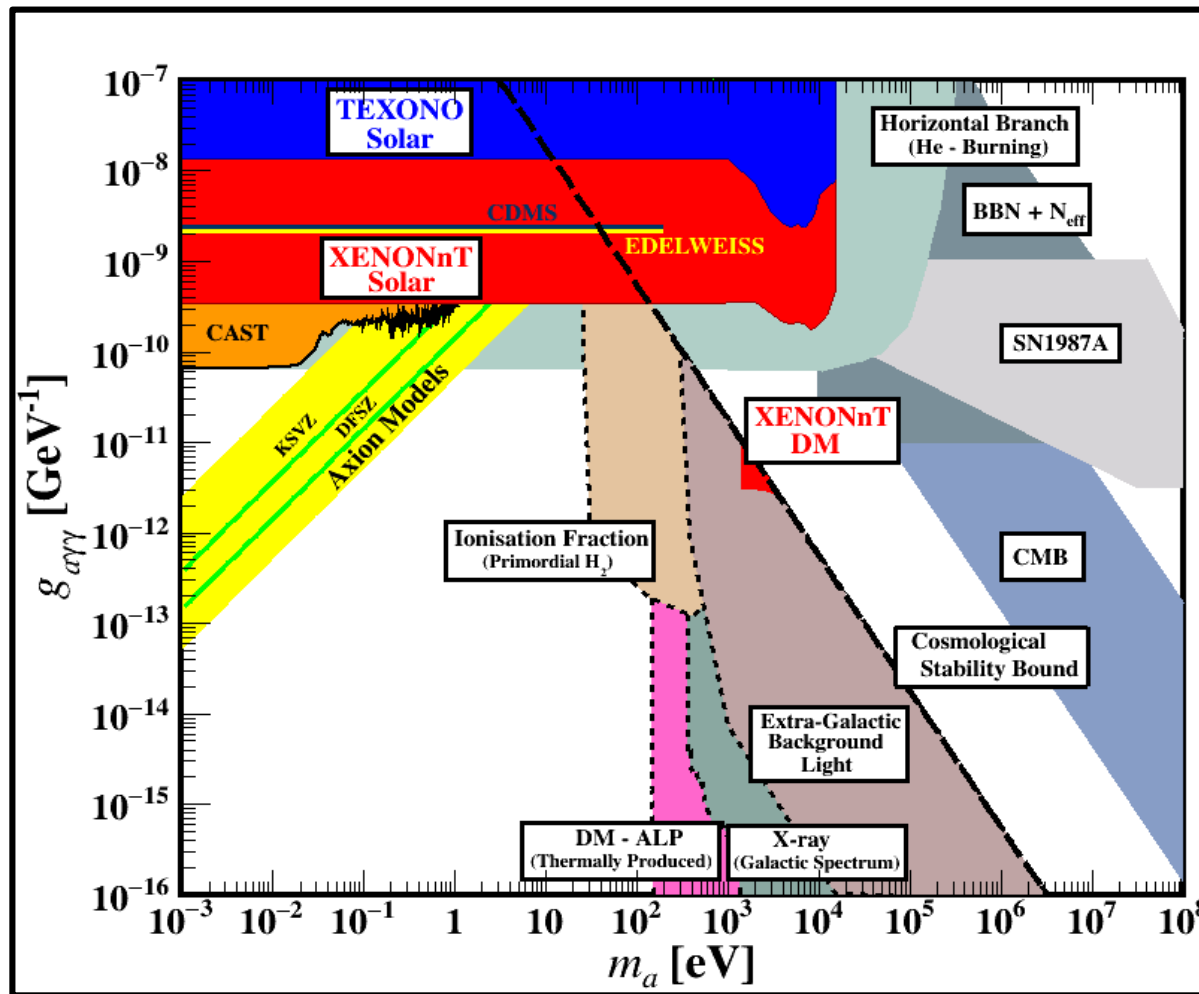
# Exclusion Plots – DM ALP



(a)



# Global Exclusion Plot

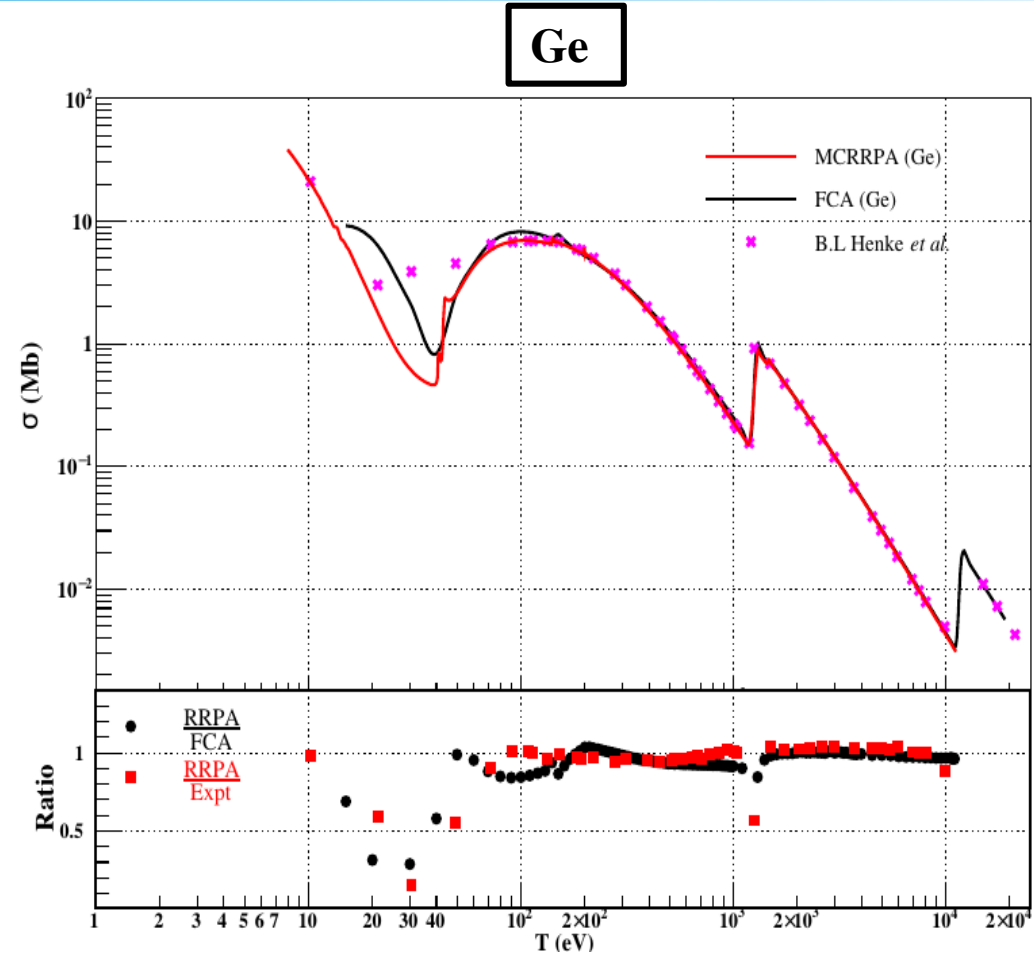
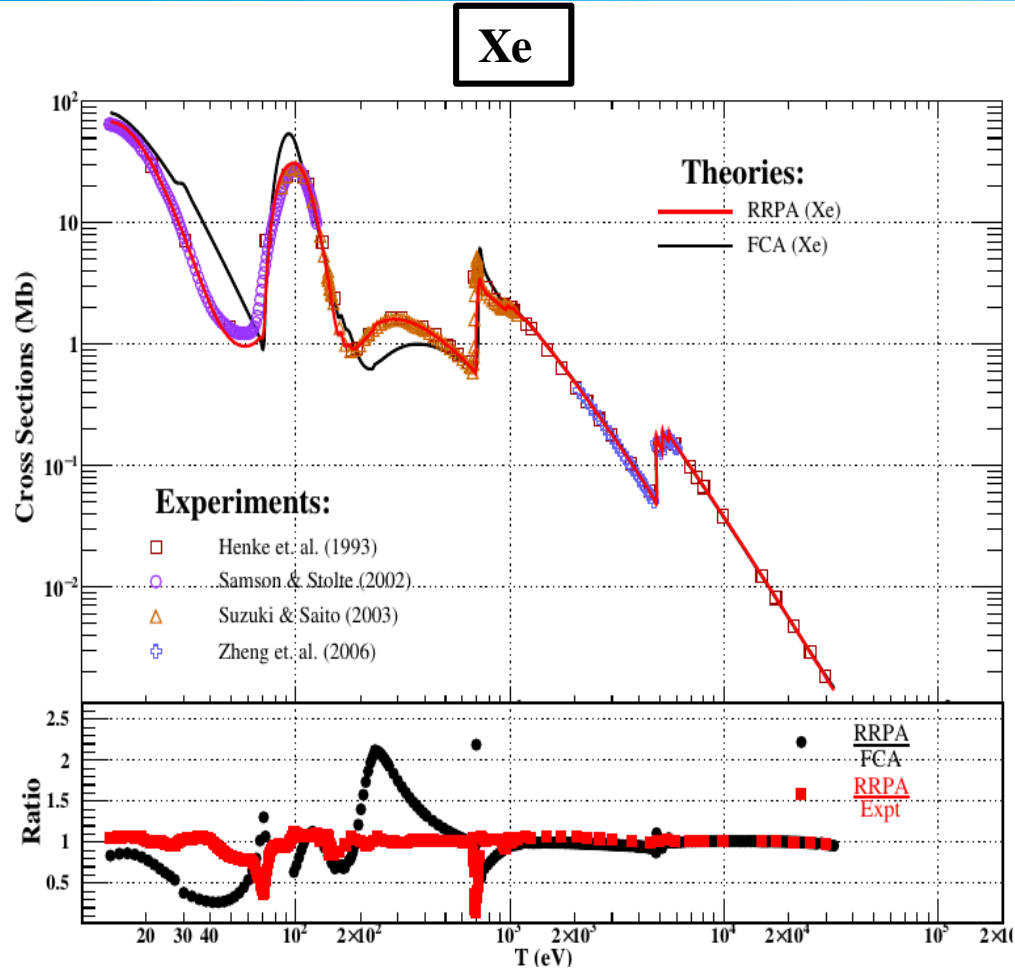


# Summary & Future Prospects

- Identified **new detection channel** ( $IP_{\text{ion}}$ ) to probe  $g_{a\gamma\gamma}$  for laboratory-based experimental searches.
- Developed **new sophisticated tools** (**FCA, RRPA, MCRRPA**) for calculating cross-section for many-body systems.
- $IP_{\text{ion}}$  channel has **discovery potential** for the search of ALPs.
- New tools are good at calculating cross-section above 80eV within **5% error**.
- Expanded the sensitivities of laboratory experiments in probing the  $m_a - g_{a\gamma\gamma}$  parameter space.
- $IP_{\text{ion}}$  sensitivities of future projects (**DARWIN**) would exceed the cosmological stability bound for DM-ALPs.
- Comprehensive analysis of ALPs through a **global approach** considering the coupling between ALPs and photons, as well as ALPs and electrons.
- Translating experimental constraints of IP effects of this work to bounds on  $g_{aee}$  based on **quantum loop effect**.

*Thank You!...*

# Photoabsorption cross section





# FCA, RRPA & MCRRPA

## ➤ Frozen Core Approximation (FCA)

- In FCA, core electrons in the system are treated as "frozen" or inert, meaning their positions are fixed and do not participate in the calculation of the  $e^- - e^-$  interactions.
- By freezing the core electrons, the  $e^- - e^-$  interactions are effectively treated as a perturbation on top of the frozen core.

## ➤ Relativistic Random Phase Approximation (RRPA)

- For many-body systems RPA can be used to account for collective effects that influence the system's response to external perturbations.

## ➤ Multi-Configuration Relativistic Random Phase Approximation (MCRRPA)

- MCRRPA is a theoretical framework used to calculate cross sections for nuclear reactions involving heavy nuclei, combining aspects of the multiconfiguration Dirac-Hartree-Fock (MCDHF) method with the RRPA to account for both relativistic effects and collective excitations.