

# RIXS Evidence for the Order–Disorder Antiferroelectric Transition in $\text{CuCrP}_2\text{S}_6$

Y. C. Su<sup>1</sup>, G. Channagowdra<sup>2</sup>, H. Y. Huang<sup>2</sup>, J. Okamoto<sup>2</sup>, A. Sudhakaran<sup>1</sup>,  
C. N. Kuo<sup>3,4</sup>, C. S. Lue<sup>3,4,5</sup>, C. T. Chen<sup>2</sup>, A. Fujimori<sup>2,1,6</sup>, C. Y. Mou<sup>1</sup> and D. J. Huang<sup>2,1,7</sup>

<sup>1</sup>Department of Physics, National Tsing Hua University, Hsinchu 30013, Taiwan

<sup>2</sup>National Synchrotron Radiation Research Center, Hsinchu 30076, Taiwan

<sup>3</sup>Department of Physics, National Cheng Kung University, Tainan 70101, Taiwan

<sup>4</sup>Taiwan Consortium of Emergent Crystalline Materials, National Science and Technology Council, Taipei 10601, Taiwan

<sup>5</sup>Program on Key Materials, Academy of Innovative Semiconductor and Sustainable Manufacturing, National Cheng Kung University, Tainan

<sup>6</sup>Department of Physics, University of Tokyo, Bunkyo-ku, Tokyo 113-0033, Japan

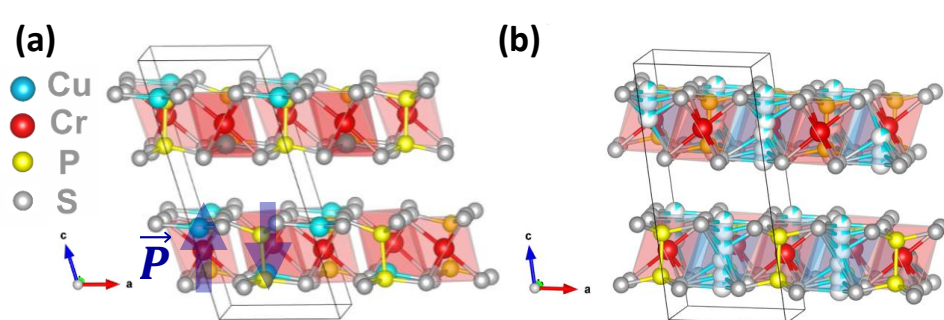
<sup>7</sup>Department of Electrophysics, National Yang Ming Chiao Tung University, Hsinchu 300093, Taiwan

## ABSTRACT

We report polarization-dependent resonant inelastic X-ray scattering (RIXS) measurements that reveal the order–disorder nature of the antiferroelectric (AFE)–paraelectric (PE) transition in the van der Waals multiferroic  $\text{CuCrP}_2\text{S}_6$ . The RIXS transition from the high-spin ground state to a state with  ${}^4T_{1g}$  symmetry splits into two  $dd$  excitations at 2.0 and 2.15 eV, arising from the splitting of  ${}^4T_{1g}$  into  ${}^4E$  and  ${}^4A_2$  due to trigonal distortion. The excitation energies remain nearly unchanged across the AFE–PE transition, indicating that the average Cr-centered trigonal crystal-field strength is insensitive to the Cu positions within RIXS resolution. However, in the AFE phase, these two peaks exhibit a distinct polarization contrast between RIXS excitations with  $\pi$  and  $\sigma$  polarizations, which evolves strongly with temperature. This polarization contrast vanishes in the PE phase, indicating that while the local trigonal splitting persists,  $\text{Cu}^+$  ions become dynamically disordered, thereby washing out the polarization anisotropy. These observations support a Cu-driven order–disorder mechanism for the AFE–PE transition, rather than a displacive transition.

## Introduction

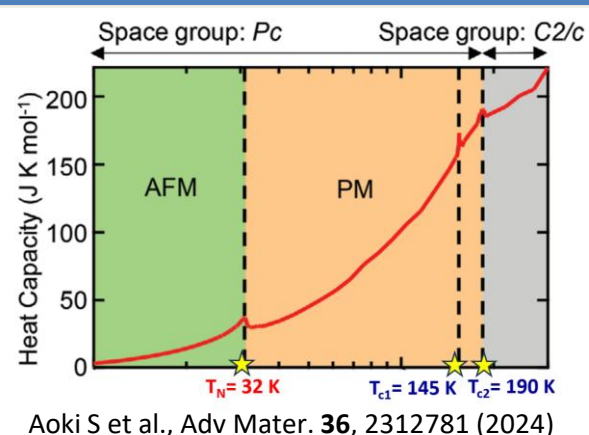
### Van der Waals multiferroic $\text{CuCrP}_2\text{S}_6$



$\text{Cu}^+$ :  $d^{10}$  Ferroelectricity  $a = 5.9098, b = 10.2447, c = 13.3644 \text{ (\AA)}$   
 $\text{Cr}^{3+}$ :  $d^3$  Ferromagnetism  $\alpha = \gamma = 90^\circ, \beta = 106.974^\circ$   
Space group: monoclinic

Above 190 K, the system crystallizes in a centrosymmetric  $C2/c$  structure, whereas below 190 K it adopts a noncentrosymmetric  $Pc$  space group. In the intermediate temperature range between 145 and 190 K, a quasi-antipolar phase emerges, followed by a fully developed antiferroelectric (AFE) phase below 145 K, characterized by long-range antipolar ordering of strongly off-centered Cu cations.

### Phase diagram

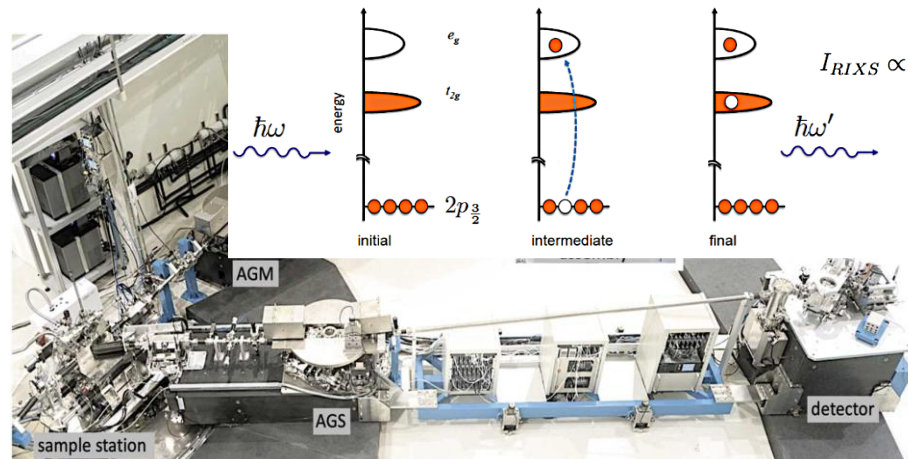


Heat-capacity measurements reveal anomalies at 32, 145, and 190 K, indicating three phase transitions. Below 32 K,  $\text{CuCrP}_2\text{S}_6$  develops an A-type antiferromagnetic order, in which Cr spins are ferromagnetically aligned within each layer and antiferromagnetically coupled between neighboring layers along the  $c$ -axis. The two higher-temperature anomalies correspond to structural transitions.

We combine temperature- and polarization-dependent resonant inelastic X-ray scattering (RIXS) to elucidate the underlying mechanisms of the AFE transition in  $\text{CuCrP}_2\text{S}_6$ . RIXS measurements probe the local electronic structure of Cr  $3d$  and its evolution throughout the transition.

## Resonant inelastic X-ray scattering (RIXS)

### TPS 41A RIXS end station



### Kramers-Heisenberg Formula

$$I_{\text{RIXS}} \propto \sum_f \left| \sum_n \frac{\langle f | \mathcal{D}_{\vec{k}, \epsilon}^1 | n \rangle \langle n | \mathcal{D}_{\vec{k}, \epsilon}^1 | i \rangle}{E_n - E_i - \hbar\omega_k + i\Gamma_n/2} \right|^2 \delta(E_f - E_i - \hbar\omega)$$

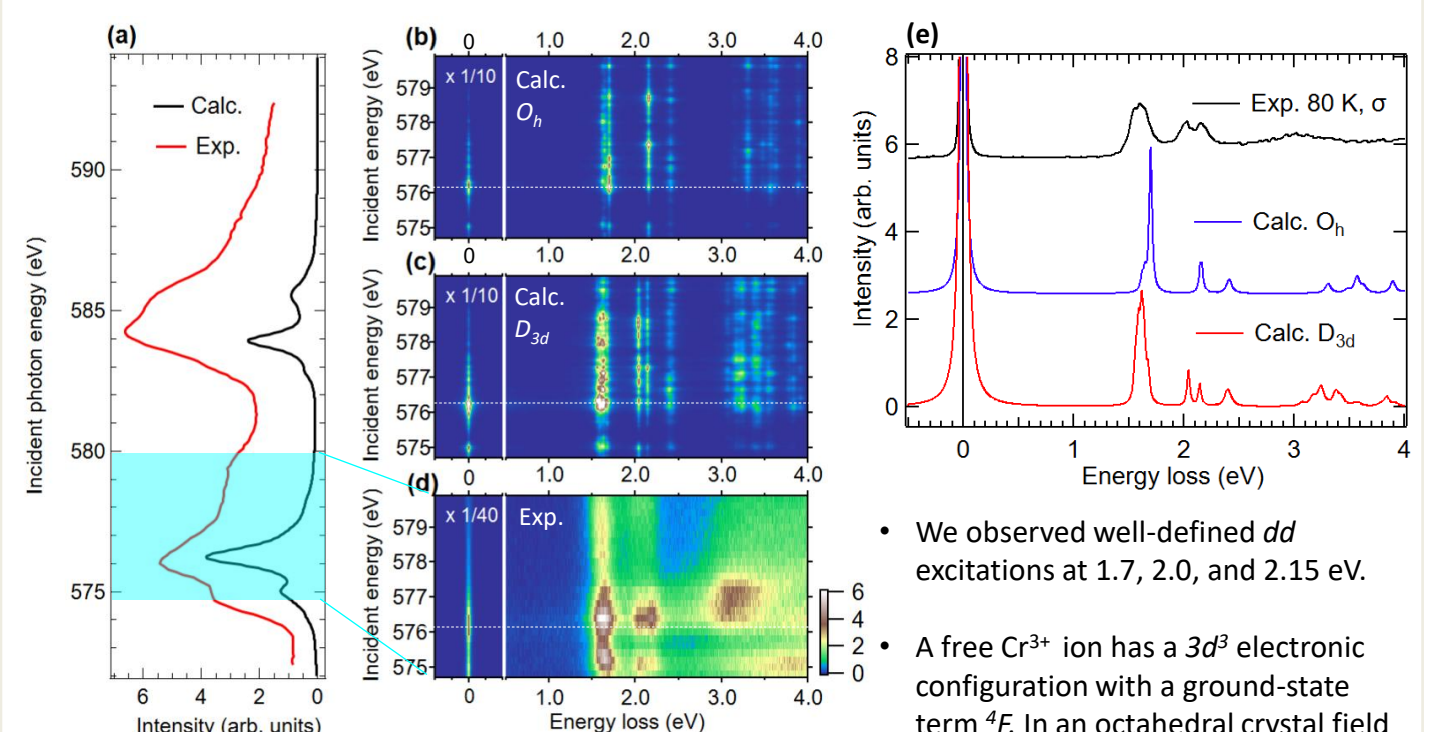
$$\mathcal{D}_{\vec{k}, \epsilon} = \sum_{i=1}^N \vec{e}_i \cdot \vec{r}_i \quad \text{Dipole operator}$$

- Two-photon process: dipole-forbidden  $dd$  excitations

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## Results & Crystal-field multiplet calculations

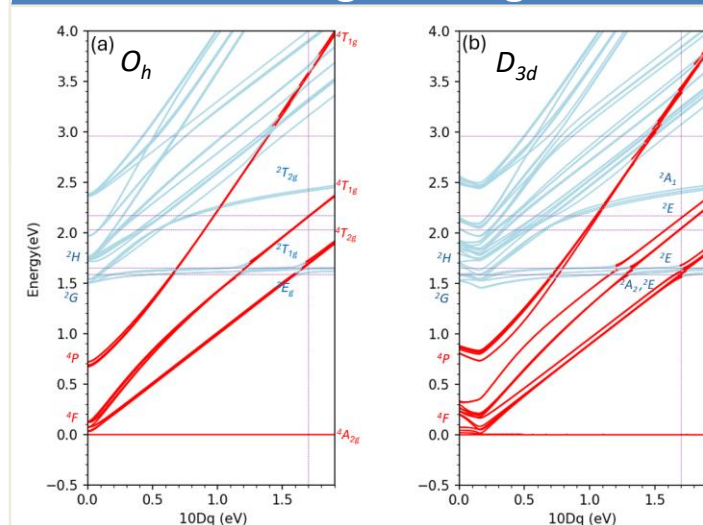
### Cr $L_{2,3}$ -edge RIXS intensity maps



Comparisons between experimental and crystal-field multiplet calculations of (a) XAS and (b)-(d) RIXS intensity maps show the crystal-field symmetry of the Cr site can be approximated by  $D_{3d}$  symmetry. The white dotted lines are the slice cuts for Cr  $L$ -edge, as shown in Fig. (e).

- We observed well-defined  $dd$  excitations at 1.7, 2.0, and 2.15 eV.
- A free  $\text{Cr}^{3+}$  ion has a  $3d^3$  electronic configuration with a ground-state term  ${}^4F$ . In an octahedral crystal field (symmetry  $O_h$ ), this term is split as:  ${}^4F \rightarrow {}^4A_{2g} + {}^4T_{2g} + {}^4T_{1g}$ , and the ground state is  ${}^4A_{2g}$ .

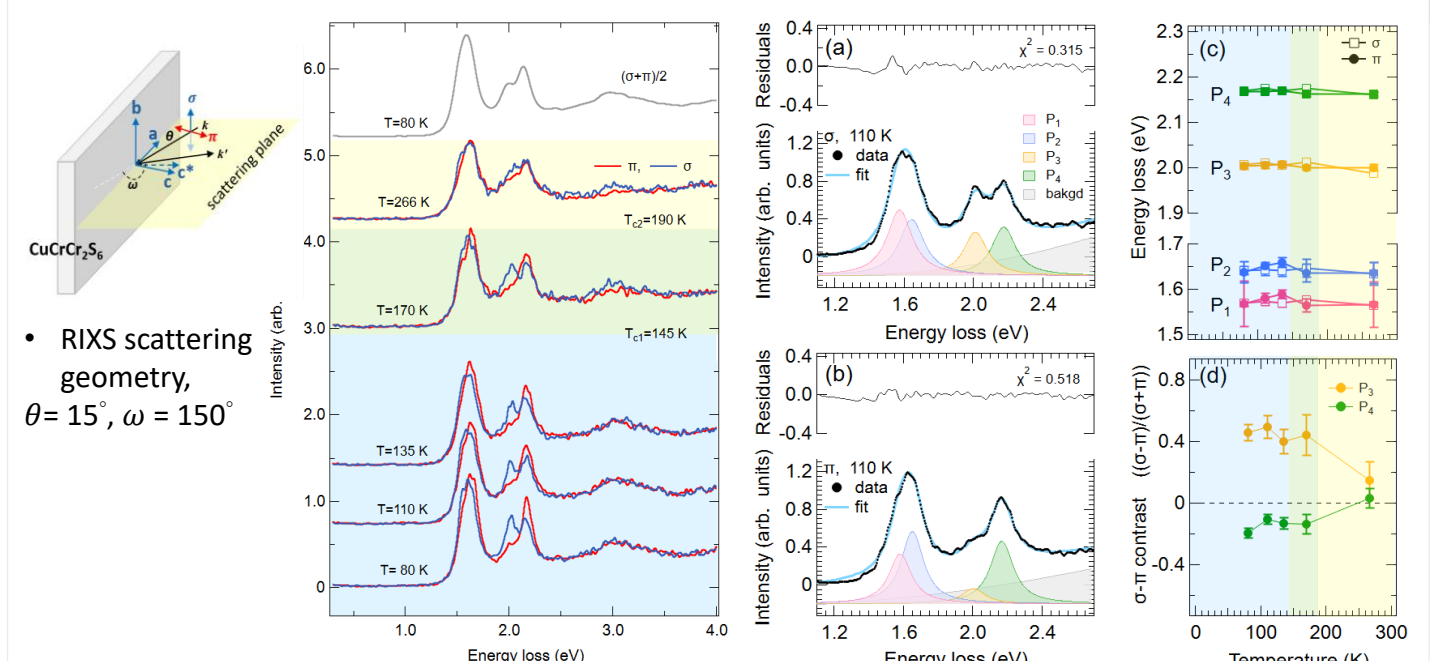
### $d^3$ Tanabe-Sugano diagram



- For a trigonal distortion corresponding to an elongation along the threefold axis, the orbitally triply degenerate  ${}^4T_{1g}$  state further splits according to  ${}^4T_{1g} \rightarrow {}^4A_2 + {}^4E$ , and the energy of the  ${}^4A_2$  state is higher than that of the  ${}^4E$  state.

- The branching of  $O_h$  to  $D_{3d}$  can be visualized by  $d^3$  Tanabe-Sugano diagram with (a)  $O_h$  and (b)  $D_{3d}$  symmetry as a function of  $10Dq$ . The purple dashed lines in the horizontal direction are the indications of multiplets observed in the RIXS spectra. The spin quartets and doublets are colored by red and blue color, respectively.

### Temperature- & Polarization-dependent RIXS



- RIXS scattering geometry,  $\theta = 15^\circ, \omega = 150^\circ$

- From the polarization-dependent RIXS results, we assign the higher-energy 2.15 eV peak to the  ${}^4A_2$  singlet and the lower-energy 2.0 eV peak to the  ${}^4E$  doublet.
- These excitation energies remain unchanged across the antiferroelectric (AFE)–paraelectric (PE) transition, indicating that the average Cr-centered trigonal crystal field is insensitive to  $\text{Cu}^+$ -driven structural changes.
- Polarization contrast between  $\pi$  and  $\sigma$  channels evolves strongly with temperature. Upon heating into the PE phase, the local symmetry axis will fluctuate away from the  $c^*$ -axis to gain entropy. The unwell-defined local symmetry makes the polarization contrast largely collapses, even though the local and instantaneous trigonal splitting remains finite.
- These findings support a Cu-driven order–disorder mechanism for the AFE–PE transition, rather than a displacive one. In the low-T AFE phase, static  $\text{Cu}^+$  ordering produces strong polarization-dependent RIXS contrast; in the PE phase, dynamic  $\text{Cu}^+$  disorder erases this anisotropy while preserving Cr-centered trigonal splitting.